# Cross-correlation between spectroscopic and photometric maps

#### Feb 13th 2014

#### **KIAS-NCTS** Joint Workshop

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#### Key observables in spectroscopic surveys

Angular diameter distance  $D_A$ : Exploiting BAO as standard rulers which measure the angular diameter distance and expansion rate as a function of redshift.

Radial distance  $H^{-1}$ : Exploiting redshift distortions as intrinsic anisotropy to decompose the radial distance represented by the inverse of Hubble rate as a function of redshift.

Coherent motion  $G_{\Theta}$ : The coherent motion, or flow, of galaxies can be statistically estimated from their effect on the clustering measurements of large redshift surveys, or through the measurement of redshift space distortions.

#### Key observables in spectroscopic surveys

Angular diameter distance  $D_A$ : Exploiting BAO as standard rulers which measure the angular diameter distance and expansion rate as a function of redshift.

Radial distance H<sup>-1</sup>: Exploiting redshift distortions as intrinsic anisoTipose are essential observables to test the inverse of Hubble rate as a function of redshift acceleration; theoretical models explaining cosmic acceleration; 1) ACDM universe

Coherent motion, Go, The coherent motion, or 9%, of galaxies can be 2) a Einstein's gravity frat large scale on the clustering measurements of large redshift surveys, or through the measurement of redshift space distortions.

# What is "Large Scale Structure formation"?

Tracing invisible ripples using pattern of light on the bottom weak lensing Tracing invisible ripples using fallen leaves on the water galaxy survey



## What is "Large Scale Structure formation"?

Tracing invisible ripples using pattern of light on the bottom weak lensing

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Signature of dark energy on LSS The same physics weakening gravity dims SN light influences on the structure of galaxy formation at large scales. without dark energy with dark energy

![](_page_5_Picture_1.jpeg)

#### Longitude problem in 21st century Measuring the redshift of the universe

![](_page_6_Picture_1.jpeg)

#### Longitude problem in 21st century Measuring the redshift of the universe

r'<17.55, d>2", 6°slice

redshift space 62295 galaxies 200 h-1 Mpc

400 h-1 Mpc

2dF scanned the first evidence of cosmic web of the universe. New spectroscopy technology allows us to locate the radial distance.

#### Measurements using large scale structure

![](_page_8_Figure_1.jpeg)

The full history of cosmic expansion can be reconstructed using galaxy redshift surveys. An anisotropy arises because of galaxy recession velocities include components from both the Hubble flow and peculiar velocities, which allows constraints to be placed on the rate of clustering growth and Hubble flow.

#### Measurements using large scale structure

![](_page_9_Picture_1.jpeg)

The full history of cosmic expansion can be reconstructed using galaxy redshift surveys. An anisotropy arises because of galaxy recession velocities include components from both the Hubble flow and peculiar velocities, which allows constraints to be placed on the rate of clustering growth and Hubble flow.

## Standard rulers of $D_A$ and $H^{-1}$

![](_page_10_Picture_1.jpeg)

![](_page_10_Picture_2.jpeg)

μ=1

![](_page_10_Picture_3.jpeg)

## Standard rulers of $D_A$ and $H^{-1}$

![](_page_11_Figure_1.jpeg)

μ=1

Blake, Galzebrook 2004 Seo, Eisenstein 2004

μ=0

#### Standard rulers of $D_A$ and $H^{-1}$

![](_page_12_Figure_1.jpeg)

μ=1

Blake, Galzebrook 2004 Seo, Eisenstein 2004

μ=0

# 2D Anisotropy Analysis

![](_page_13_Figure_1.jpeg)

# 2D Anisotropy Analysis

![](_page_14_Figure_1.jpeg)

# 2D Anisotropy Analysis

![](_page_15_Figure_1.jpeg)

## Can we formulate RD in precision cosmology?

Finger of God effect at small scales

(Jackson 1972)

Squeezing effect at large scales

(Kaiser 1987)

$$\begin{split} & \bigvee \\ \mathsf{P}_{\mathsf{s}}(\mathsf{k},\mu) = [\mathsf{P}_{\mathsf{gg}}(\mathsf{k}) + \Delta \mathsf{P}_{\mathsf{gg}} + 2\mu^2 \mathsf{P}_{\mathsf{g}\Theta}(\mathsf{k}) + \Delta \mathsf{P}_{\mathsf{g}\Theta} + \mu^4 \mathsf{P}_{\Theta\Theta}(\mathsf{k}) + \Delta \mathsf{P}_{\Theta\Theta} \\ & \quad + \mu^2 \mathsf{A}(\mathsf{k}) + \mu^4 \mathsf{B}(\mathsf{k}) + \mu^6 \mathsf{C}(\mathsf{k}) + \dots ] \exp[-(\mathsf{k}\mu\sigma_{\mathsf{p}})^2] \end{split}$$

 $P_{s}(k,\mu) = P_{qq}(k) + 2\mu^{2}P_{q\Theta}(k) + \mu^{4}P_{\Theta\Theta}(k)$ 

Taruya, Nichimishi, Saito 2010

#### Conversion into configuration space $P_{s}(k,\mu) = [Q_{0}(k) + \mu^{2}Q_{2}(k) + \mu^{4}Q_{4}(k) + \mu^{6}Q_{6}(k)] \exp[-(k\mu\sigma_{p})^{2}]$

YSS, Okumura, Taruya 2013 Taruya, Nichimishi, Saito 2010

Conversion into configuration space  $P_{s}(k,\mu) = [Q_{0}(k) + \mu^{2}Q_{2}(k) + \mu^{4}Q_{4}(k) + \mu^{6}Q_{6}(k)] \exp[-(k\mu\sigma_{p})^{2}]$  $\xi(\sigma,\pi) = \int d^3 k P(k,\mu) e^{ikx} = \Sigma \xi_{\ell}(s) \mathcal{P}_{\ell}(v)$  $\xi_{\ell}(s) = i^{\ell} \int k^2 dk P_{\ell}(k) j_{\ell}(ks)$  $P_0(k) = p_0(k)$  $P_2(k) = 5/2 [3p_1(k) - p_0(k)]$  $P_4(k) = 9/8 [35p_2(k) - 30p_1(k) + 3p_0(k)]$  $P_6(k) = \frac{13}{16} [231p_3(k) - 315p_2(k) - 105p_1(k) + 5p_0(k)]$  $p_n(k) = 1/2 \left[ \frac{\gamma(n+1/2,\kappa)}{\kappa^{n+1/2}Q_0(k)} + \frac{\gamma(n+3/2,\kappa)}{\kappa^{n+3/2}Q_2(k)} \right]$ +  $\gamma(n+5/2,\kappa)/\kappa^{n+5/2}Q_4(k) + \gamma(n+7/2,\kappa)/\kappa^{n+7/2}Q_6(k)$  $K = k^2 \sigma^2_p$ 

YSS, Okumura, Taruya 2013 Taruya, Nichimishi, Saito 2010

# Theoretical reproduction of 2D BAO $\xi(\sigma,\pi)$ with Kaiser+Gaussian FoG

![](_page_19_Figure_1.jpeg)

YSS, Okumura, Taruya 2013

# Theoretical reproduction of 2D BAO $\xi(\sigma,\pi)$ with additional correction terms

![](_page_20_Figure_1.jpeg)

YSS, Okumura, Taruya 2013

#### **Observed spectra in Fourier space** $P_{s}(k,\mu) = P_{gg}(k) + 2\mu^{2} \sqrt{P_{gg}(k)} P_{\Theta\Theta}(k) + \mu^{4} P_{\Theta\Theta}(k)$

![](_page_21_Figure_1.jpeg)

#### **Observed spectra in Fourier space** $P_{s}(k,\mu) = P_{gg}(k) + 2\mu^{2}\sqrt{P_{gg}(k)}P_{\Theta\Theta}(k) + \mu^{4}P_{\Theta\Theta}(k)$

![](_page_22_Figure_1.jpeg)

 $\mu$ 

# Effect of non-linear correction terms Theoretical model of RSD is broken down at $k_{max}=0.1h/Mpc$ .

![](_page_23_Figure_1.jpeg)

YSS, Nichimishi, Taruya, Kayo 2013

#### Dark Energy Spectroscopy Instrument

![](_page_24_Picture_1.jpeg)

#### Dark Energy Probes

Providing 50M LRG or ELG galaxies at northern hemisphere, to observe BAO and RSD which are essential for dark energy.

#### Probing Neutrino mass

Observing high redshift quasars to get Lyman alpha forest, and providing linear spectra at small scaes to probe neutrino mass.

#### KASI's participation from 2013

KASI will become an associate institute from 2013. We take part in scientific efforts to probe dark energy using spectroscopy wide-deep field survey. We start to invite researchers to collaborate DESI project.

#### RD version of Alcock-Paczynski effect

![](_page_25_Figure_1.jpeg)

YSS 2013

The maximal information from RD Maximal Fisher matrix description in 3D  $F_{\alpha\beta} = \int \frac{d^{3}k}{2(2\pi)^{3}} \frac{\partial P(k,\mu)}{\partial p_{\alpha}} \frac{\partial P(k,\mu)}{[P(k,\mu) + 1/n_{gal}]^{2}} \frac{\partial P(k,\mu)}{\partial p_{\beta}}$   $\sim \frac{1}{2(2\pi)^{3}} \sum \sum_{\mu=0}^{\mu=1} \sum_{\mu=1}^{\mu=1} \sum_{\mu=0}^{\mu=1} \sum_{\mu=1}^{\mu=1} \sum_$ 

YSS, Gaztanaga, Martin prepared

The alternate method using 2D The expression can be effectively reduced using 2D projected spectra

$$F_{\alpha\beta} = \Sigma \frac{\partial C_{ij}}{\partial p_{\alpha}} \Theta^{-1}_{ij,mn} \frac{\partial C_{mn}}{\partial p_{\beta}}$$

We probe both radial and transverse distances, and coherent motions along the line of sight using spectroscopy surveys. But when we use the 2D projected spectra, it is not quite clear how we can include the information of radial distance and coherent motion growth function. Solution: thin slicing approximation for C<sub>l</sub> Full AP effect and the influence of coherent motion can be employed using thin enough binning

$$F_{\alpha\beta} = \Sigma \frac{\partial C_{ij}}{\partial p_{\alpha}} \Theta^{-1}_{ij,mn} \frac{\partial C_{mn}}{\partial p_{\beta}}$$

sical length

; Y H Q

correlation length

Bin #2

Bin #1

Therefore cross correlation between two bins are not fully independent

Gaztanaga et.al. 2011

The agreement is under investigation Maximal Fisher matrix description in 3D  $F_{\alpha\beta} = \int \frac{d^{3}k}{2(2\pi)^{3}} \frac{\partial P(k,\mu)}{\partial p_{\alpha}} \frac{\partial P(k,\mu)}{[P(k,\mu) + 1/n_{aal}]^{2}} \frac{\partial P(k,\mu)}{\partial p_{\beta}}$  $\sim \frac{1}{2(2\pi)^3} \sum_{\mu=0}^{3} \sum_{\mu=0}^{\mu=1} |^2 \Delta_{\mu} \Delta_{\mu} \frac{\partial P(l,\mu)}{\partial p_{\alpha}} \frac{\partial P(l,\mu)}{[P(l,\mu) + 1/n_{ad}]^2} \frac{\partial P(l,\mu)}{\partial p_{\beta}}$ Full AP effect and the influence of coherent motion can be employed using thin enough binning  $F_{\alpha\beta} = \Sigma \frac{\partial C_{ij}}{\partial p_{\alpha}} \Theta^{-1}_{ij,mn} \frac{\partial C_{mn}}{\partial p_{\beta}}$ 

YSS, Gaztanaga, Martin prepared

The agreement is under investigation Maximal Fisher matrix description in 3D  $F_{\alpha\beta} = \int \frac{d^{3}k}{2(2\pi)^{3}} \frac{\partial P(k,\mu)}{\partial p_{\alpha}} \frac{\partial P(k,\mu)}{[P(k,\mu) + 1/n_{aal}]^{2}} \frac{\partial P(k,\mu)}{\partial p_{\beta}}$  $\sim \frac{1}{2(2\pi)^3} \sum_{\mu=0}^{3} \sum_{\mu=0}^{3} |2\Delta|\Delta\mu| \frac{\partial P(l,\mu)}{\partial p_{\alpha}} \frac{\partial P(l,\mu)}{[P(l,\mu) + 1/n_{gal}]^2} \frac{\partial P(l,\mu)}{\partial p_{\beta}}$ Full AP effect and the influence of coherent motion can be employed using thin enough binning Due to some unknown reasons, the consistency is 

YSS, Gaztanaga, Martin prepared

#### The benefit of 2D expression

more comprehensive in cross correlation between spectroscopy and photometry maps

$$F_{\alpha\beta} = \Sigma \frac{\partial C_{ij}}{\partial p_{\alpha}} \Theta^{-1}_{ij,mn} \frac{\partial C_{mn}}{\partial p_{\beta}}$$

#### The benefit of 2D expression

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$$F_{\alpha\beta} = \Sigma \frac{\partial C_{ij}}{\partial p_{\alpha}} \Theta^{-1}_{ij,mn} \frac{\partial C_{mn}}{\partial p_{\beta}}$$

![](_page_32_Figure_3.jpeg)

#### The benefit of 2D expression

more comprehensive in cross correlation between spectroscopy and photometry maps

$$F_{\alpha\beta} = \Sigma \frac{\partial C_{ij}}{\partial p_{\alpha}} \Theta^{-1}_{ij,mn} \frac{\partial C_{mn}}{\partial p_{\beta}}$$

![](_page_33_Figure_3.jpeg)

![](_page_34_Figure_0.jpeg)

#### Implication of cosmic acceleration

Breaking down our knowledge of particle physics: we have limited knowledge of particle physics bounded by testable high energy, and our efforts to explain the cosmic acceleration turn out in vain.

Alternative mechanism to generate fine tuned vacuum energy

New unknown energy component

Unification or coupling between dark sectors

Breaking down our knowledge of gravitational physics: gravitational physics has been tested in solar system scales, and it is yet confirmed at horizon size,

Presence of extra dimension

Non-linear interaction to Einstein equation

Failure of standard cosmology model: our understanding of the universe is still standing on assumptions:

Inhomogeneous models: LTB, back reaction

## Motivation for the same sky survey

![](_page_36_Figure_1.jpeg)

#### Coupling between dark sectors Baryon CDM

![](_page_37_Picture_1.jpeg)

 $Q_c^{\mu} = -\alpha \rho_c \nabla^{\mu} \varphi \,,$ 

If the coupling term is proportional to scalar field, then Euler equation is broken, i.e. the universality of free falling between baryon and dark matter is violated.

#### Brane world model (DGP)

Brane (X=0)

In DGP, gravity alone propagates in the bulk, and 5D gravitational theory is complemented by an induced 4D Ricci scalar restricted to the brane.

 $S = \int d^{5}x \sqrt{-g} [R^{(5)}/2\mu^{(5)} + \delta(\chi)(R^{(4)}/2\mu^{(4)} + L_{M})]$ 

#### Minkowski bulk

The ratio of the two scales defines a cross-over radius beyond which the four dimensional gravitational theory transition into a five dimensional regime,

 $r_c = \mu^{(5)}/2\mu^{(4)}$ 

Dvali, Gabadadze, Porrati (2000)

#### Brane world model (DGP)

Sound wave propagate through air (leaving) table, and system on the table loses energy at  $r>r_c$ .

Although balls do not leave, energy loss due to leaking sound wane above table.

rc

# f(R) gravity

Corrections are introduced in the Einstein-Hilbert Lagrangian to modify the general relativity, which gets influential only low curvature, e.g. late time & not dense region. The corrections can be adjusted to generate the cosmic acceleration, Carroll, Duvvuri, Trodden, Turner (2004:CDTT)

$$S = \int d^4x \sqrt{-g} \left[ rac{R+f(R)}{2\mu^2} + \mathcal{L}_{
m m} 
ight]$$

cosmic acceleration was discovered with  $f(R) = \frac{1}{2}a/R$ . Ruled out

Two distinct branches of f(R) gravity was found depending on the sign of second order derivative of f(R) in terms of R,

> $f_{RR} = d^2 f/dR^2 < 0$  Unstable  $f_{RR} = d^2 f/dR^2 > 0$  Stable

The original proposal of CDTT is ruled out due to instability.

YSS, Hu, Sawicki (2007)

## Motivation for the same sky survey

![](_page_41_Figure_1.jpeg)

#### Motivation for the same sky survey

![](_page_42_Figure_1.jpeg)

#### Structure formation

![](_page_43_Figure_1.jpeg)

#### What is good to scan both at the same fields?

photometry  $\Phi-\Psi$ 

spectroscopy

δ

At different regions, we will make only parametrized test on different dark energy models.

#### What is good to scan both at the same fields?

spleotoosebpy ΦδΨ

With scanning the same fields, we are able to make fully covariant test of different dark energy models.

#### Test of mass screening effect

#### $C_{l}^{KK} = \Sigma \Delta D W^2 \Delta^2_{\phi\phi}$

![](_page_46_Picture_2.jpeg)

![](_page_46_Picture_3.jpeg)

Song, Kazuya 2010

# Test of mass screening effect

 $C_{l}^{KK} = \Sigma \Delta D W^{2} \Delta^{2}_{\phi\phi}$  $C_{l}^{KK} = \Sigma \Delta D W^{'2} C_{l}^{\delta\delta}$ 

Dynamical DE  $C_{l}^{KK} = C_{l}^{KK}$ 

 $\Delta^2_{\Phi\Phi} C_{\Gamma}^{\delta\delta}$ 

Geometrical DE  $C_{l}^{KK} \neq C_{l}^{KK}$ 

![](_page_47_Picture_4.jpeg)

 $k^2 \Phi = 4 \pi G_N a^2 \rho_m \delta_m$ 

Song, Kazuya 2010

#### Consistency test using cross-correlation

$$E_G(R) = \frac{1}{\beta} \frac{\Upsilon_{\rm gm}(R)}{\Upsilon_{\rm gg}(R)}.$$

![](_page_48_Figure_2.jpeg)

Zhang et.al. 2009 Reyes et.al. 2010

#### Generation of wide field survey

#### **Photometric (WL)** Spectroscopic (BAO, topology, coherent motions)

	Ground	Space
Second 2004-2009	CFTHLS	
	SDSS II —	→ 0.1 <z<0.5, n="100,000&lt;/td"></z<0.5,>
Third 2011-2015	DES	
	SDSS III (BOSS)—	→ 0.2 <z<0.8, n="1,500,000&lt;/td"></z<0.8,>
Next 2018-2023	LSST	WFIRST, EUCLID
	DESI	WFIRST, EUCLID

#### Missing component at ground telescope

![](_page_50_Figure_1.jpeg)

## Missing component at ground telescope

![](_page_51_Figure_1.jpeg)

#### 4m Multi Object Spectroscopic Telescope

#### Major Science Drivers

#### GAIA follow up

Providing full 6D coordinates for stellar science in Milky way. Aiming dynamical structure of Milky Way.

#### eROSITA follow up

Surveying >50,000 southern galaxy clusters. Investigating the evolution of galaxy populations in clusters.

#### 4m Multi Object Spectroscopic Telescope

#### Major Science Drivers

#### GAIA follow up

Providing full 6D coordinates for stellar science in Milky way. Aiming dynamical structure of Milky Way.

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Surveying >50,000 southern galaxy clusters. Investigating the evolution of galaxy populations in clusters.

#### Alternative Science Goals

4MOST will provide a unique opportunity to determine the nature of dark energy by cross-correlating photometric and spectroscopic wide deep field surveys in southern sky.

# European spectroscopy wide deep survey in Chile

![](_page_54_Picture_1.jpeg)

#### THE window for the same sky survey

	Baseline	Goal
class	4m	-
FoV	3 deg	>5 deg
Fiber pos	1500	>3000
Area	10000deg <sup>2</sup>	20000deg <sup>2</sup>
Objects	5M	>20M

#### European spectroscopy wide deep survey

4MO

# The ultimate cosmological test of Einstein's gravity

![](_page_55_Picture_2.jpeg)

#### Concepts of participation

4MOS

![](_page_56_Picture_1.jpeg)

#### AAO-KASI collaboration

# $w_{L} measures \Phi - \Psi$ $w_{L$

#### Conclusion

We are interested in the same sky survey of the photometry and spectroscopy wide deep survey in southern hemisphere. Enhancing dark energy constraint by the same sky survey Key observables to test general relativity cosmologically Understanding the nature of dark energy in model independent way

We plan to launch the collaboration group to study strong lens probe of the universe.

> The follow-up strong lensing observation targeted by DES or LSST Probing dark matter and dark energy simultaneously An independent probe for dark energy science

We are interested in starting asteroid search using 0.5m telescope in southern hemisphere.